

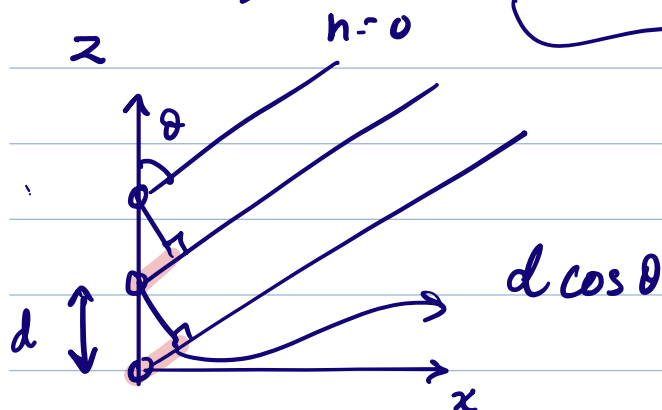
Antenna arrays & Fourier theory

We now explore the connections between antenna array theory & signal processing.

① Start with the array factor (AF) for a uniformly spaced 1D array (spacing d , wavevector k):

We know $E_{\text{total}}(\theta) = f(\theta) \text{AF}(\theta)$ (by pattern multiplication)

$$\text{AF}(\theta) = \sum_{n=0}^{N-1} a_n e^{jk d \cos \theta n}$$



complex array excitation (mag) and (phase)

Set $\psi = k d \cos \theta$

$$\text{AF}(\psi) = \sum_{n=0}^{N-1} a_n e^{jn\psi}$$

② Let's analyse $\text{AF}(\psi)$. Is it periodic?

$\text{AF}(\psi + 2\pi) = \text{AF}(\psi)$. Yes, with period 2π , for any $\{a_n\}$. \Rightarrow We only need to know AF over one period, e.g. $[-\pi, \pi)$.

③ A common statement is "the AF is just a Fourier transform". Is it true?

Consider 3 transforms:

<u>Transform</u>	<u>Defn</u>	<u>Argument</u>
1) DTFT	$A(\psi) = \sum_{n=0}^{N-1} a_n e^{jn\psi}$	ψ conts.
2) DFT	$A[m] = \sum_{n=0}^{N-1} a_n e^{j\frac{2\pi}{N}nm}$	m discrete $m \in [0, N-1]$
3) z-transform	$A(z) = \sum_{n=0}^{N-1} a_n z^{-n}$	$z \in \mathbb{C}$

aside

Reminder of ①: $A(\omega) \triangleq \int_{-\infty}^{\infty} a(t) e^{-j\omega t} dt$

Sample it every T : $A(\omega) \Big|_{\omega/T} = \sum_{-\infty}^{\infty} \underbrace{a(nT)T}_{a_n} e^{-j\omega nT}$

$= \sum_{-\infty}^{\infty} a_n e^{-jn(\omega T)}$

So, what is our AF? DTFT of $\{a_n\}$ evaluated at a contns frequency ψ .

Then if we sample the DTFT at discrete freqs $\psi = \frac{2\pi m}{N}$, we get the DFT of $\{a_n\}$

i.e. $A[m] = AF\left(\psi = \frac{2\pi m}{N}\right)$.

So if we do $AF_samples = fft(a, M)$
what are we doing? $(M \log M)$

1) M -point FFT

$\Rightarrow M$ uniform samples in $[0, 2\pi)$
i.e. $\psi = [0, \frac{2\pi}{N}, \frac{2\pi}{N} \cdot 2 \dots, \frac{2\pi}{N}(M-1)]$

2) If $M > N$?

$\{a_n\}$ is zero padded to length M

Summary: The AF is technically a DTFT,
when we take a FFT of $\{a_n\}$ we are sampling
the DTFT.

④ The visible region.

As $\theta \rightarrow 0$ to π ,

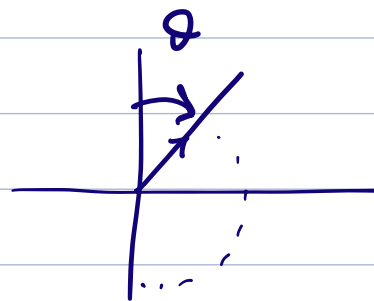
$u = \cos \theta$ goes

$u \rightarrow 1$ to -1

and $\psi = kd \cos \theta$ goes

$\psi \rightarrow kd$ to $-kd$

i.e. $\theta \rightarrow [0, \pi]$ maps $\psi [-kd, kd]$



In ψ space, we are spanning $kd + kd$
 $= 2kd \rightarrow$ defined as the visible
region since it corresponds
to $\theta = 0$ to π (physically
visible)

[Recall this is a 1D problem, symmetric about z axis, so no need for $\theta > \pi$]

$$\text{Now: } 2kd = 2 \cdot \frac{2\pi}{\lambda} \cdot d = \frac{4\pi d}{\lambda}$$

and the periodicity of AF is 2π .

$$\hookrightarrow \text{At } d = \lambda/2, 2kd = 2\pi$$

\Rightarrow Every feature of AF(ψ) is seen once.

$$\hookrightarrow \text{if } d < \lambda/2? \quad 2kd < 2\pi$$

\Rightarrow As θ goes from $0 \rightarrow \pi$, AF plot covers less than 2π .

\Rightarrow Some features of AF are hidden.

$$\hookrightarrow \text{if } d > \lambda/2? \quad 2kd > 2\pi$$

\Rightarrow As θ goes from $0 \rightarrow \pi$, AF plot covers more than 2π .

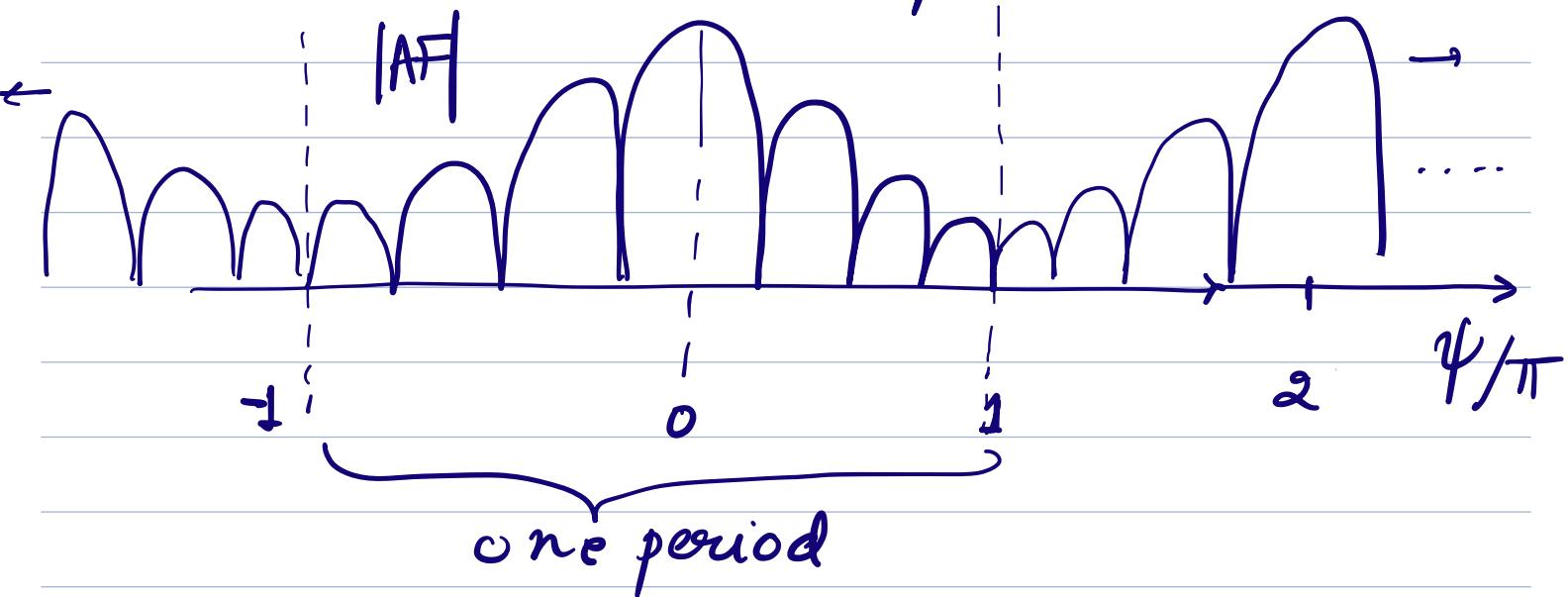
\rightarrow But AF is still periodic with 2π .

\Rightarrow Parts of adjacent periods enter the visible region \rightarrow Grating lobes (an alternate explanation).

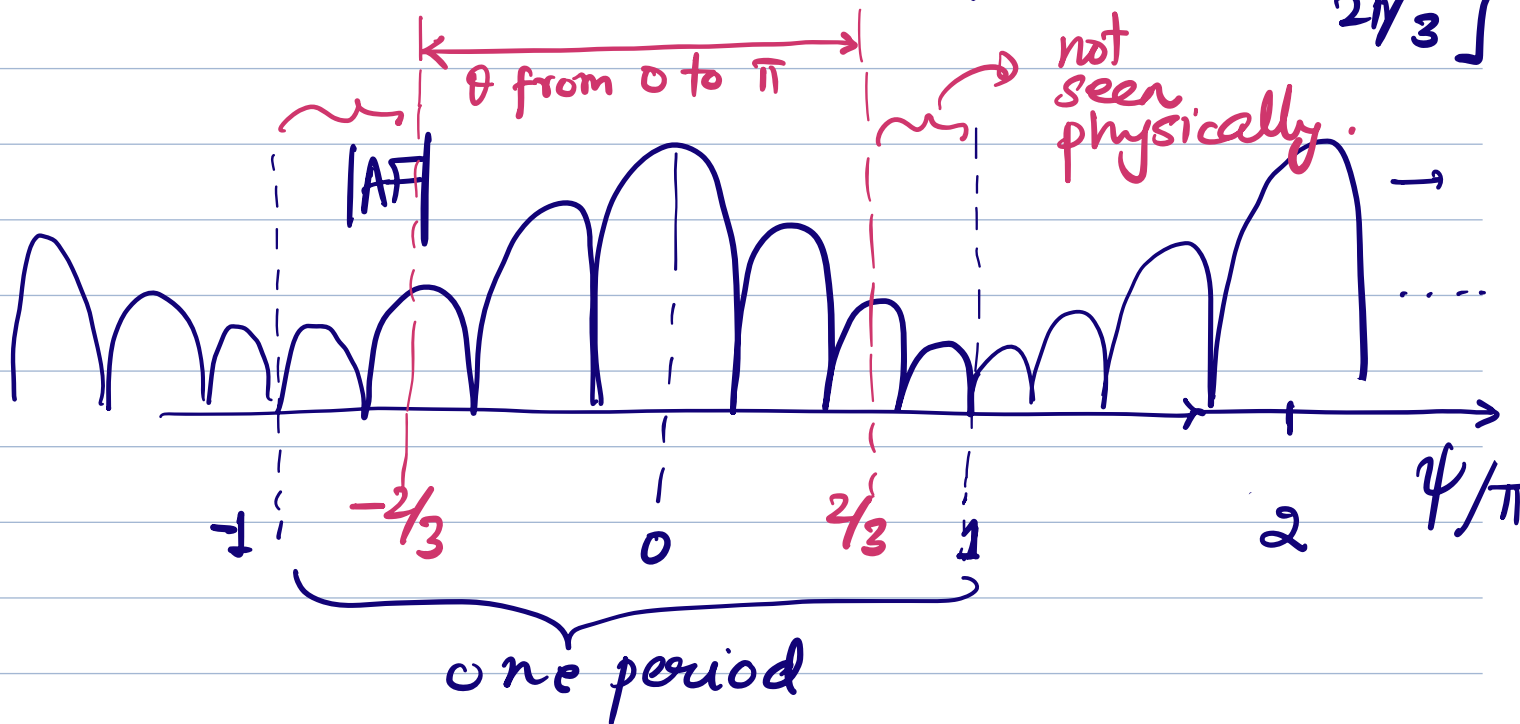
Numerical example with $N=8$, $a_n = 1 \forall n$.

$$AF = \frac{\sin N\psi/2}{\sin \psi/2}, \quad \psi = kd \cos \theta$$

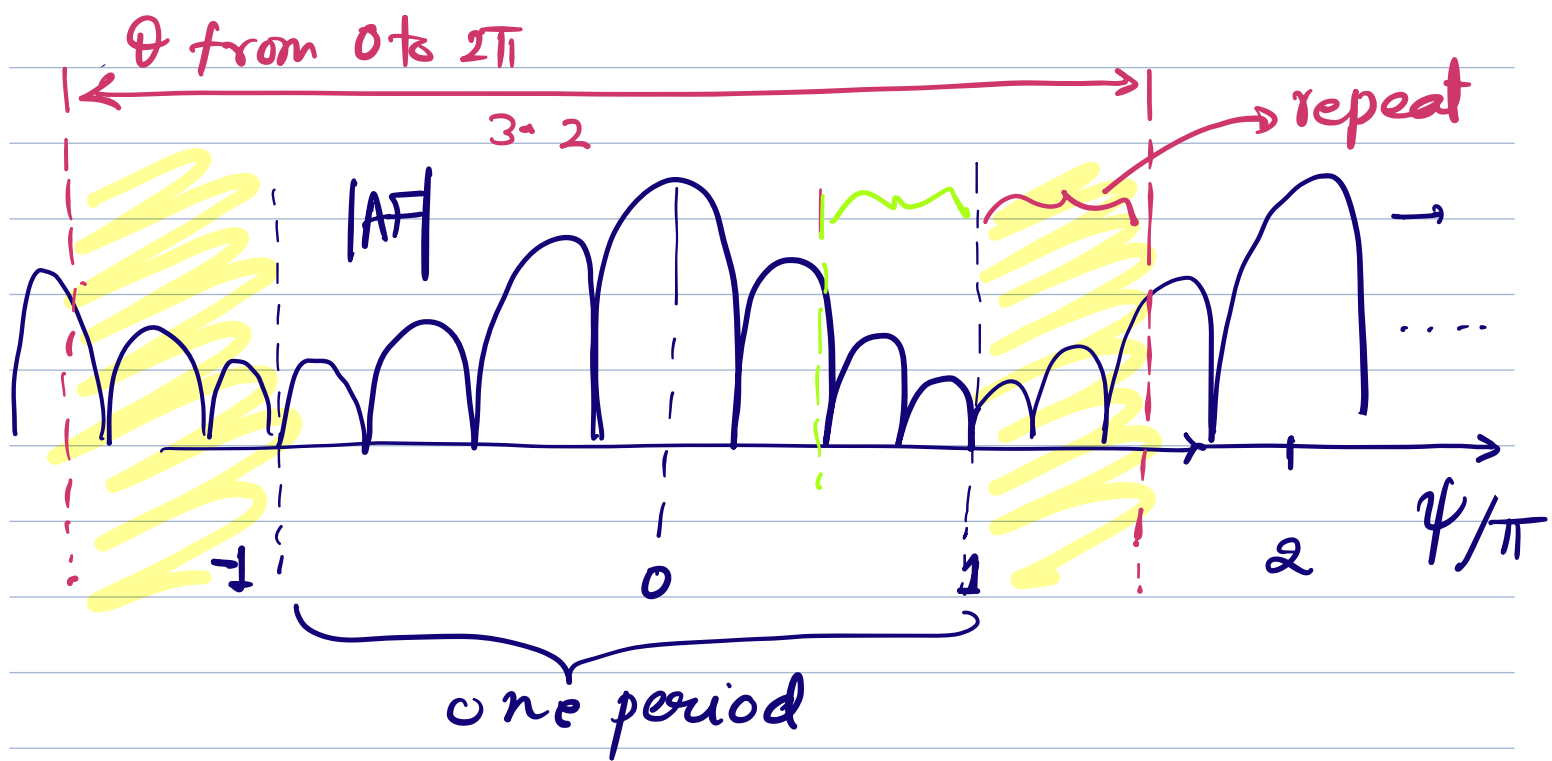
\therefore AF vs ψ doesn't depend on kd .



\hookrightarrow case 1, $d = \lambda/3 \Rightarrow kd = 2\pi/3$, visible $\psi \in [-2\pi/3, 2\pi/3]$



\hookrightarrow case 2, $d = 0.8\lambda \Rightarrow kd = 1.6\pi$
 \Rightarrow visible region $\psi \in [-1.6\pi, 1.6\pi]$

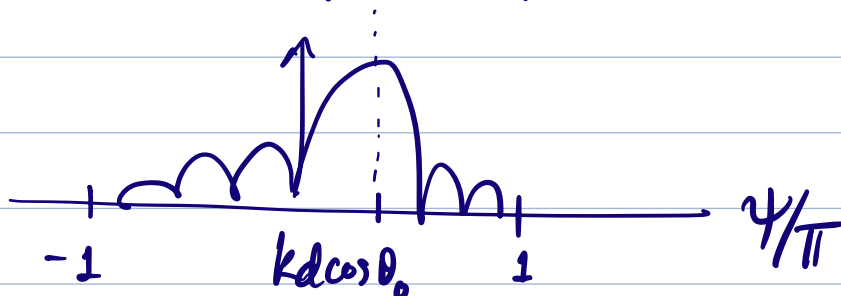


The shaded regions are being seen twice when θ goes from 0 to π .

When $d = 1\lambda$, $kd = 4\pi \Rightarrow$ we see the main lobe from other periods enter into the visible region.
Mostly undesirable!

A note: if progressive phase shift, then $AF = \sum a_n e^{jn\psi'}$
with $\psi' = kd \cos\theta + \beta$. If we want beam to form at $\theta = \theta_0$, then $AF = \frac{\sin N\psi'/2}{\sin \psi'/2} \Rightarrow kd \cos\theta_0 + \beta = 0$.

with $\psi' = \psi + \beta$, the $AF(\psi)$ plot will have a maxima at $\psi = -\beta = kd \cos\theta_0$



↳ In summary, analogy with the sampling theorem.

Signal $f(t)$,
conts time,
with b/w B

→ sampled every $t = T_s$ ($f_s = \frac{1}{T_s}$)



Nyquist condition, $f_s > 2B$
else, overlapping (aliasing)

In the antenna array case?

Aperture distribution
 $a(z)$ (continuous z)

→ sampled every
 $z = d$, i.e. $f_s = \frac{1}{d}$.



Spatial Nyquist condn
if $f_s = \frac{1}{d} > \frac{2}{\lambda}$ (i.e. $d < \frac{\lambda}{2}$)

else, grating lobes

⇒ We can think of the signal $a(z)$ having
spatial bandwidth = $\frac{1}{\lambda}$

(by analogy with B
above)